

Problem 4.37

- (a) Apply S_- to $|10\rangle$ (Equation 4.175), and confirm that you get $\sqrt{2}\hbar|1-1\rangle$.
- (b) Apply S_{\pm} to $|00\rangle$ (Equation 4.176), and confirm that you get zero.
- (c) Show that $|11\rangle$ and $|1-1\rangle$ (Equation 4.175) are eigenstates of S^2 , with the appropriate eigenvalue.

Solution**Part (a)**

Equation 4.175 is on page 177. Use the spin relationships in Equation 4.146 on page 168.

$$\begin{aligned}
 S_-|10\rangle &= S_- \left[\frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle) \right] & (4.175) \\
 &= \frac{1}{\sqrt{2}} \left[S_-|\uparrow\downarrow\rangle + S_-|\downarrow\uparrow\rangle \right] \\
 &= \frac{1}{\sqrt{2}} \left[(S_-^{(1)} + S_-^{(2)})|\uparrow\downarrow\rangle + (S_-^{(1)} + S_-^{(2)})|\downarrow\uparrow\rangle \right] \\
 &= \frac{1}{\sqrt{2}} \left[(S_-^{(1)}|\uparrow\rangle)|\downarrow\rangle + |\uparrow\rangle(S_-^{(2)}|\downarrow\rangle) \right. \\
 &\quad \left. + (S_-^{(1)}|\downarrow\rangle)|\uparrow\rangle + |\downarrow\rangle(S_-^{(2)}|\uparrow\rangle) \right] \\
 &= \frac{1}{\sqrt{2}} \left[(\hbar|\downarrow\rangle)|\downarrow\rangle + |\uparrow\rangle(\mathbf{0}) + (\mathbf{0})|\uparrow\rangle + |\downarrow\rangle(\hbar|\downarrow\rangle) \right] \\
 &= \frac{1}{\sqrt{2}} (\hbar|\downarrow\downarrow\rangle + \hbar|\downarrow\downarrow\rangle) \\
 &= \frac{1}{\sqrt{2}} (2\hbar|\downarrow\downarrow\rangle) \\
 &= \sqrt{2}\hbar|\downarrow\downarrow\rangle \\
 &= \sqrt{2}\hbar|1-1\rangle
 \end{aligned}$$

Part (b)

Equation 4.176 is on page 177 too. Use the spin relationships in Equation 4.146 on page 168.

$$\begin{aligned}
 S_+|00\rangle &= S_+ \left[\frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle) \right] & S_-|00\rangle &= S_- \left[\frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle) \right] \\
 &= \frac{1}{\sqrt{2}} \left[S_+|\uparrow\downarrow\rangle - S_+|\downarrow\uparrow\rangle \right] & &= \frac{1}{\sqrt{2}} \left[S_-|\uparrow\downarrow\rangle - S_-|\downarrow\uparrow\rangle \right] \\
 &= \frac{1}{\sqrt{2}} \left[(S_+^{(1)} + S_+^{(2)})|\uparrow\downarrow\rangle - (S_+^{(1)} + S_+^{(2)})|\downarrow\uparrow\rangle \right] & &= \frac{1}{\sqrt{2}} \left[(S_-^{(1)} + S_-^{(2)})|\uparrow\downarrow\rangle - (S_-^{(1)} + S_-^{(2)})|\downarrow\uparrow\rangle \right] \\
 &= \frac{1}{\sqrt{2}} \left[(S_+^{(1)}|\uparrow\rangle)|\downarrow\rangle + |\uparrow\rangle(S_+^{(2)}|\downarrow\rangle) \right. & &= \frac{1}{\sqrt{2}} \left[(S_-^{(1)}|\uparrow\rangle)|\downarrow\rangle + |\uparrow\rangle(S_-^{(2)}|\downarrow\rangle) \right. \\
 &\quad \left. - (S_+^{(1)}|\downarrow\rangle)|\uparrow\rangle - |\downarrow\rangle(S_+^{(2)}|\uparrow\rangle) \right] & &\quad \left. - (S_-^{(1)}|\downarrow\rangle)|\uparrow\rangle - |\downarrow\rangle(S_-^{(2)}|\uparrow\rangle) \right] \\
 &= \frac{1}{\sqrt{2}} \left[(\mathbf{0})|\downarrow\rangle + |\uparrow\rangle(\hbar|\uparrow\rangle) \right. & &= \frac{1}{\sqrt{2}} \left[(\hbar|\downarrow\rangle)|\downarrow\rangle + |\uparrow\rangle(\mathbf{0}) \right. \\
 &\quad \left. - (\hbar|\uparrow\rangle)|\uparrow\rangle - |\downarrow\rangle(\mathbf{0}) \right] & &\quad \left. - (\mathbf{0})|\uparrow\rangle - |\downarrow\rangle(\hbar|\downarrow\rangle) \right] \\
 &= \frac{1}{\sqrt{2}} (\hbar|\uparrow\uparrow\rangle - \hbar|\uparrow\uparrow\rangle) & &= \frac{1}{\sqrt{2}} (\hbar|\downarrow\downarrow\rangle - \hbar|\downarrow\downarrow\rangle) \\
 &= \mathbf{0} & &= \mathbf{0}
 \end{aligned}$$

Part (c)

Let S^2 act on $|11\rangle$ first. Use the relationships in Equation 4.171 on page 176.

$$\begin{aligned}
S^2|11\rangle &= (\mathbf{S} \cdot \mathbf{S})|\uparrow\uparrow\rangle \\
&= (\mathbf{S}^{(1)} + \mathbf{S}^{(2)}) \cdot (\mathbf{S}^{(1)} + \mathbf{S}^{(2)})|\uparrow\uparrow\rangle \\
&= \left(\mathbf{S}^{(1)} \cdot \mathbf{S}^{(1)} + \mathbf{S}^{(1)} \cdot \mathbf{S}^{(2)} + \mathbf{S}^{(2)} \cdot \mathbf{S}^{(1)} + \mathbf{S}^{(2)} \cdot \mathbf{S}^{(2)} \right) |\uparrow\uparrow\rangle \\
&= \left[(S^{(1)})^2 + 2\mathbf{S}^{(1)} \cdot \mathbf{S}^{(2)} + (S^{(2)})^2 \right] |\uparrow\uparrow\rangle \\
&= (S^{(1)})^2 |\uparrow\uparrow\rangle + (S^{(2)})^2 |\uparrow\uparrow\rangle + 2\mathbf{S}^{(1)} \cdot \mathbf{S}^{(2)} |\uparrow\uparrow\rangle \\
&= (S^{(1)})^2 \left| \frac{1}{2} \frac{1}{2} \frac{1}{2} \frac{1}{2} \right\rangle + (S^{(2)})^2 \left| \frac{1}{2} \frac{1}{2} \frac{1}{2} \frac{1}{2} \right\rangle + 2\mathbf{S}^{(1)} \cdot \mathbf{S}^{(2)} |\uparrow\uparrow\rangle \\
&= \frac{1}{2} \left(\frac{1}{2} + 1 \right) \hbar^2 \left| \frac{1}{2} \frac{1}{2} \frac{1}{2} \frac{1}{2} \right\rangle + \frac{1}{2} \left(\frac{1}{2} + 1 \right) \hbar^2 \left| \frac{1}{2} \frac{1}{2} \frac{1}{2} \frac{1}{2} \right\rangle + 2\mathbf{S}^{(1)} \cdot \mathbf{S}^{(2)} |\uparrow\uparrow\rangle \\
&= \frac{3\hbar^2}{2} \left| \frac{1}{2} \frac{1}{2} \frac{1}{2} \frac{1}{2} \right\rangle + 2\mathbf{S}^{(1)} \cdot \mathbf{S}^{(2)} |\uparrow\uparrow\rangle \\
&= \frac{3\hbar^2}{2} |\uparrow\uparrow\rangle + 2 \left\langle S_x^{(1)}, S_y^{(1)}, S_z^{(1)} \right\rangle \cdot \left\langle S_x^{(2)}, S_y^{(2)}, S_z^{(2)} \right\rangle |\uparrow\uparrow\rangle \\
&= \frac{3\hbar^2}{2} |\uparrow\uparrow\rangle + 2 \left\langle \frac{1}{2}(S_+^{(1)} + S_-^{(1)}), \frac{1}{2i}(S_+^{(1)} - S_-^{(1)}), S_z^{(1)} \right\rangle \cdot \left\langle \frac{1}{2}(S_+^{(2)} + S_-^{(2)}), \frac{1}{2i}(S_+^{(2)} - S_-^{(2)}), S_z^{(2)} \right\rangle |\uparrow\uparrow\rangle \\
&= \frac{3\hbar^2}{2} |\uparrow\uparrow\rangle + \left[\frac{1}{2}(S_+^{(1)} + S_-^{(1)})(S_+^{(2)} + S_-^{(2)}) - \frac{1}{2}(S_+^{(1)} - S_-^{(1)})(S_+^{(2)} - S_-^{(2)}) + 2S_z^{(1)}S_z^{(2)} \right] |\uparrow\uparrow\rangle \\
&= \frac{3\hbar^2}{2} |\uparrow\uparrow\rangle + \frac{1}{2}(S_+^{(1)} + S_-^{(1)})(S_+^{(2)} + S_-^{(2)}) |\uparrow\uparrow\rangle - \frac{1}{2}(S_+^{(1)} - S_-^{(1)})(S_+^{(2)} - S_-^{(2)}) |\uparrow\uparrow\rangle + 2S_z^{(1)}S_z^{(2)} |\uparrow\uparrow\rangle \\
&= \frac{3\hbar^2}{2} |\uparrow\uparrow\rangle + \frac{1}{2}(S_+^{(1)} + S_-^{(1)})(S_+^{(2)} |\uparrow\uparrow\rangle + S_-^{(2)} |\uparrow\uparrow\rangle) - \frac{1}{2}(S_+^{(1)} - S_-^{(1)})(S_+^{(2)} |\uparrow\uparrow\rangle - S_-^{(2)} |\uparrow\uparrow\rangle) \\
&\hspace{20em} + 2S_z^{(1)} \left[|\uparrow\uparrow\rangle (S_z^{(2)} |\uparrow\uparrow\rangle) \right] \\
&= \frac{3\hbar^2}{2} |\uparrow\uparrow\rangle + \frac{1}{2} \left(\cancel{S_+^{(1)} S_+^{(2)} |\uparrow\uparrow\rangle} + S_+^{(1)} S_-^{(2)} |\uparrow\uparrow\rangle + S_-^{(1)} S_+^{(2)} |\uparrow\uparrow\rangle + \cancel{S_-^{(1)} S_-^{(2)} |\uparrow\uparrow\rangle} \right) \\
&\quad - \frac{1}{2} \left(\cancel{S_+^{(1)} S_+^{(2)} |\uparrow\uparrow\rangle} - S_+^{(1)} S_-^{(2)} |\uparrow\uparrow\rangle - S_-^{(1)} S_+^{(2)} |\uparrow\uparrow\rangle + \cancel{S_-^{(1)} S_-^{(2)} |\uparrow\uparrow\rangle} \right) + 2S_z^{(1)} \left[|\uparrow\uparrow\rangle \left(\frac{\hbar}{2} |\uparrow\uparrow\rangle \right) \right]
\end{aligned}$$

Continue the simplification.

$$\begin{aligned}
 S^2|11\rangle &= \frac{3\hbar^2}{2}|\uparrow\uparrow\rangle + S_+^{(1)}S_-^{(2)}|\uparrow\uparrow\rangle + S_-^{(1)}S_+^{(2)}|\uparrow\uparrow\rangle + \hbar S_z^{(1)}|\uparrow\uparrow\rangle \\
 &= \frac{3\hbar^2}{2}|\uparrow\uparrow\rangle + S_+^{(1)}\left[|\uparrow\rangle\left(S_-^{(2)}|\uparrow\rangle\right)\right] + S_-^{(1)}\left[|\uparrow\rangle\left(S_+^{(2)}|\uparrow\rangle\right)\right] + \hbar\left(S_z^{(1)}|\uparrow\rangle\right)|\uparrow\rangle \\
 &= \frac{3\hbar^2}{2}|\uparrow\uparrow\rangle + S_+^{(1)}\left[|\uparrow\rangle\left(\hbar|\downarrow\rangle\right)\right] + S_-^{(1)}\left[|\uparrow\rangle\left(\mathbf{0}\right)\right] + \hbar\left(\frac{\hbar}{2}|\uparrow\rangle\right)|\uparrow\rangle \\
 &= \frac{3\hbar^2}{2}|\uparrow\uparrow\rangle + \hbar S_+^{(1)}|\uparrow\downarrow\rangle + \frac{\hbar^2}{2}|\uparrow\uparrow\rangle \\
 &= 2\hbar^2|\uparrow\uparrow\rangle + \hbar\left(S_+^{(1)}|\uparrow\rangle\right)|\downarrow\rangle \\
 &= 2\hbar^2|\uparrow\uparrow\rangle + \hbar(\mathbf{0})|\downarrow\rangle \\
 &= 2\hbar^2|\uparrow\uparrow\rangle \\
 &= 2\hbar^2|11\rangle
 \end{aligned}$$

Now let S^2 act on $|1-1\rangle$.

$$\begin{aligned}
 S^2|1-1\rangle &= (\mathbf{S} \cdot \mathbf{S})|\downarrow\downarrow\rangle \\
 &= (\mathbf{S}^{(1)} + \mathbf{S}^{(2)}) \cdot (\mathbf{S}^{(1)} + \mathbf{S}^{(2)})|\downarrow\downarrow\rangle \\
 &= \left(\mathbf{S}^{(1)} \cdot \mathbf{S}^{(1)} + \mathbf{S}^{(1)} \cdot \mathbf{S}^{(2)} + \mathbf{S}^{(2)} \cdot \mathbf{S}^{(1)} + \mathbf{S}^{(2)} \cdot \mathbf{S}^{(2)}\right)|\downarrow\downarrow\rangle \\
 &= \left[\left(S^{(1)}\right)^2 + 2\mathbf{S}^{(1)} \cdot \mathbf{S}^{(2)} + \left(S^{(2)}\right)^2\right]|\downarrow\downarrow\rangle \\
 &= \left(S^{(1)}\right)^2|\downarrow\downarrow\rangle + \left(S^{(2)}\right)^2|\downarrow\downarrow\rangle + 2\mathbf{S}^{(1)} \cdot \mathbf{S}^{(2)}|\downarrow\downarrow\rangle \\
 &= \left(S^{(1)}\right)^2\left|\frac{1}{2}\frac{1}{2}\frac{-1}{2}\frac{-1}{2}\right\rangle + \left(S^{(2)}\right)^2\left|\frac{1}{2}\frac{1}{2}\frac{-1}{2}\frac{-1}{2}\right\rangle + 2\mathbf{S}^{(1)} \cdot \mathbf{S}^{(2)}|\downarrow\downarrow\rangle \\
 &= \frac{1}{2}\left(\frac{1}{2} + 1\right)\hbar^2\left|\frac{1}{2}\frac{1}{2}\frac{-1}{2}\frac{-1}{2}\right\rangle + \frac{1}{2}\left(\frac{1}{2} + 1\right)\hbar^2\left|\frac{1}{2}\frac{1}{2}\frac{-1}{2}\frac{-1}{2}\right\rangle + 2\mathbf{S}^{(1)} \cdot \mathbf{S}^{(2)}|\downarrow\downarrow\rangle \\
 &= \frac{3\hbar^2}{2}\left|\frac{1}{2}\frac{1}{2}\frac{-1}{2}\frac{-1}{2}\right\rangle + 2\mathbf{S}^{(1)} \cdot \mathbf{S}^{(2)}|\downarrow\downarrow\rangle \\
 &= \frac{3\hbar^2}{2}|\downarrow\downarrow\rangle + 2\left\langle S_x^{(1)}, S_y^{(1)}, S_z^{(1)} \right\rangle \cdot \left\langle S_x^{(2)}, S_y^{(2)}, S_z^{(2)} \right\rangle|\downarrow\downarrow\rangle
 \end{aligned}$$

Continue the simplification, noting that $S_x = \frac{1}{2}(S_+ + S_-)$ and $S_y = \frac{1}{2i}(S_+ - S_-)$ (page 168).

$$\begin{aligned}
S^2|1 -1\rangle &= \frac{3\hbar^2}{2}|\downarrow\downarrow\rangle + 2\left\langle\frac{1}{2}(S_+^{(1)} + S_-^{(1)}), \frac{1}{2i}(S_+^{(1)} - S_-^{(1)}), S_z^{(1)}\right\rangle \cdot \left\langle\frac{1}{2}(S_+^{(2)} + S_-^{(2)}), \frac{1}{2i}(S_+^{(2)} - S_-^{(2)}), S_z^{(2)}\right\rangle|\downarrow\downarrow\rangle \\
&= \frac{3\hbar^2}{2}|\downarrow\downarrow\rangle + \left[\frac{1}{2}(S_+^{(1)} + S_-^{(1)})(S_+^{(2)} + S_-^{(2)}) - \frac{1}{2}(S_+^{(1)} - S_-^{(1)})(S_+^{(2)} - S_-^{(2)}) + 2S_z^{(1)}S_z^{(2)}\right]|\downarrow\downarrow\rangle \\
&= \frac{3\hbar^2}{2}|\downarrow\downarrow\rangle + \frac{1}{2}(S_+^{(1)} + S_-^{(1)})(S_+^{(2)} + S_-^{(2)})|\downarrow\downarrow\rangle - \frac{1}{2}(S_+^{(1)} - S_-^{(1)})(S_+^{(2)} - S_-^{(2)})|\downarrow\downarrow\rangle + 2S_z^{(1)}S_z^{(2)}|\downarrow\downarrow\rangle \\
&= \frac{3\hbar^2}{2}|\downarrow\downarrow\rangle + \frac{1}{2}(S_+^{(1)} + S_-^{(1)})(S_+^{(2)}|\downarrow\downarrow\rangle + S_-^{(2)}|\downarrow\downarrow\rangle) - \frac{1}{2}(S_+^{(1)} - S_-^{(1)})(S_+^{(2)}|\downarrow\downarrow\rangle - S_-^{(2)}|\downarrow\downarrow\rangle) \\
&\qquad\qquad\qquad + 2S_z^{(1)}\left[|\downarrow\downarrow\rangle\left(S_z^{(2)}|\downarrow\downarrow\rangle\right)\right] \\
&= \frac{3\hbar^2}{2}|\downarrow\downarrow\rangle + \frac{1}{2}\left(\cancel{S_+^{(1)}S_+^{(2)}|\downarrow\downarrow\rangle} + S_+^{(1)}S_-^{(2)}|\downarrow\downarrow\rangle + S_-^{(1)}S_+^{(2)}|\downarrow\downarrow\rangle + \cancel{S_-^{(1)}S_-^{(2)}|\downarrow\downarrow\rangle}\right) \\
&\quad - \frac{1}{2}\left(\cancel{S_+^{(1)}S_+^{(2)}|\downarrow\downarrow\rangle} - S_+^{(1)}S_-^{(2)}|\downarrow\downarrow\rangle - S_-^{(1)}S_+^{(2)}|\downarrow\downarrow\rangle + \cancel{S_-^{(1)}S_-^{(2)}|\downarrow\downarrow\rangle}\right) + 2S_z^{(1)}\left[|\downarrow\downarrow\rangle\left(-\frac{\hbar}{2}|\downarrow\downarrow\rangle\right)\right] \\
&= \frac{3\hbar^2}{2}|\downarrow\downarrow\rangle + S_+^{(1)}S_-^{(2)}|\downarrow\downarrow\rangle + S_-^{(1)}S_+^{(2)}|\downarrow\downarrow\rangle - \hbar S_z^{(1)}|\downarrow\downarrow\rangle \\
&= \frac{3\hbar^2}{2}|\downarrow\downarrow\rangle + S_+^{(1)}\left[|\downarrow\downarrow\rangle\left(S_-^{(2)}|\downarrow\downarrow\rangle\right)\right] + S_-^{(1)}\left[|\downarrow\downarrow\rangle\left(S_+^{(2)}|\downarrow\downarrow\rangle\right)\right] - \hbar\left(S_z^{(1)}|\downarrow\downarrow\rangle\right)|\downarrow\downarrow\rangle \\
&= \frac{3\hbar^2}{2}|\downarrow\downarrow\rangle + S_+^{(1)}\left[|\downarrow\downarrow\rangle(\mathbf{0})\right] + S_-^{(1)}\left[|\downarrow\downarrow\rangle(\hbar|\uparrow\rangle)\right] - \hbar\left(-\frac{\hbar}{2}|\downarrow\downarrow\rangle\right)|\downarrow\downarrow\rangle \\
&= \frac{3\hbar^2}{2}|\downarrow\downarrow\rangle + \hbar S_-^{(1)}|\downarrow\uparrow\rangle + \frac{\hbar^2}{2}|\downarrow\downarrow\rangle \\
&= 2\hbar^2|\downarrow\downarrow\rangle + \hbar\left(S_-^{(1)}|\downarrow\downarrow\rangle\right)|\uparrow\rangle \\
&= 2\hbar^2|\downarrow\downarrow\rangle + \hbar(\mathbf{0})|\uparrow\rangle \\
&= 2\hbar^2|\downarrow\downarrow\rangle \\
&= 2\hbar^2|1 -1\rangle
\end{aligned}$$

Therefore, $|11\rangle$ and $|1-1\rangle$ are eigenstates of S^2 , and the eigenvalue of S^2 is $2\hbar^2$.